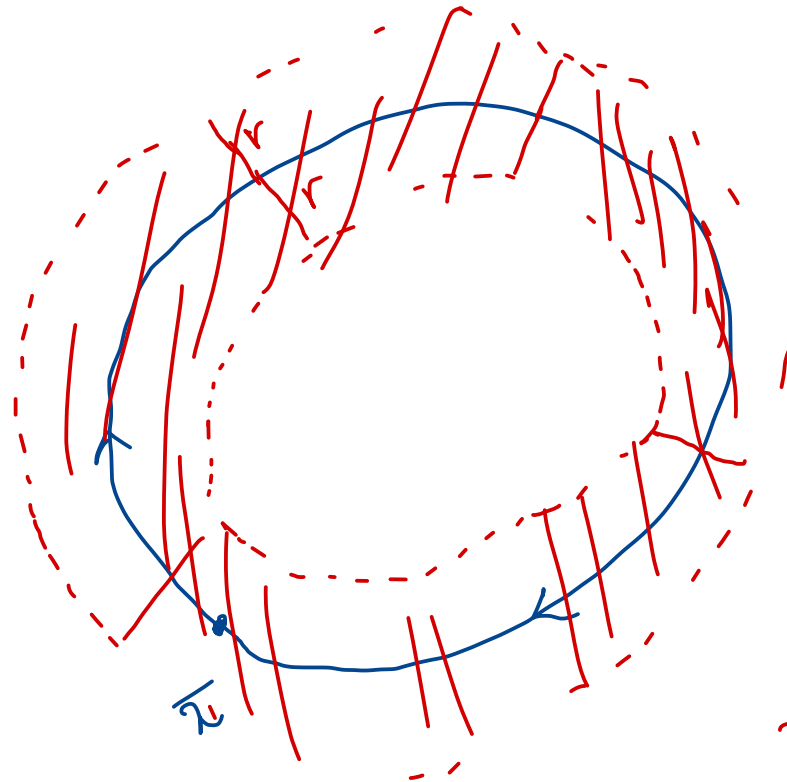


5.4. Periodic orbits: stability, limit cycles and Poincaré maps

Definition 5.31. Consider a system $\dot{x} = f(x)$, where $f : \mathbb{R}^n \rightarrow \mathbb{R}^n$ is smooth. Let $\bar{x} \in \mathbb{R}^n$ and assume it is not an equilibrium point. Then, $\Gamma = \{\phi(t, \bar{x}), t \in \mathbb{R}\}$ is said to be a periodic orbit or a cycle if there exists $T > 0$, such that $\phi(T, \bar{x}) = \bar{x}$.

- Geometrically, a periodic orbit is a closed curve.
- Assume Γ is a periodic orbit and consider $r > 0$. Let $U_r(\Gamma)$ be the set of all points $x \in \mathbb{R}^n$ whose distance from Γ is less than r .

$U_r(\Gamma)$



$T > 0$

$\phi(T, \bar{x}) = \bar{x}$

$\phi(T+t, \bar{x}) = \phi(t, \bar{x})$

$\forall t \in \mathbb{R}$

$x \in \mathbb{R}^n$

$dist(x, \Gamma) < r$

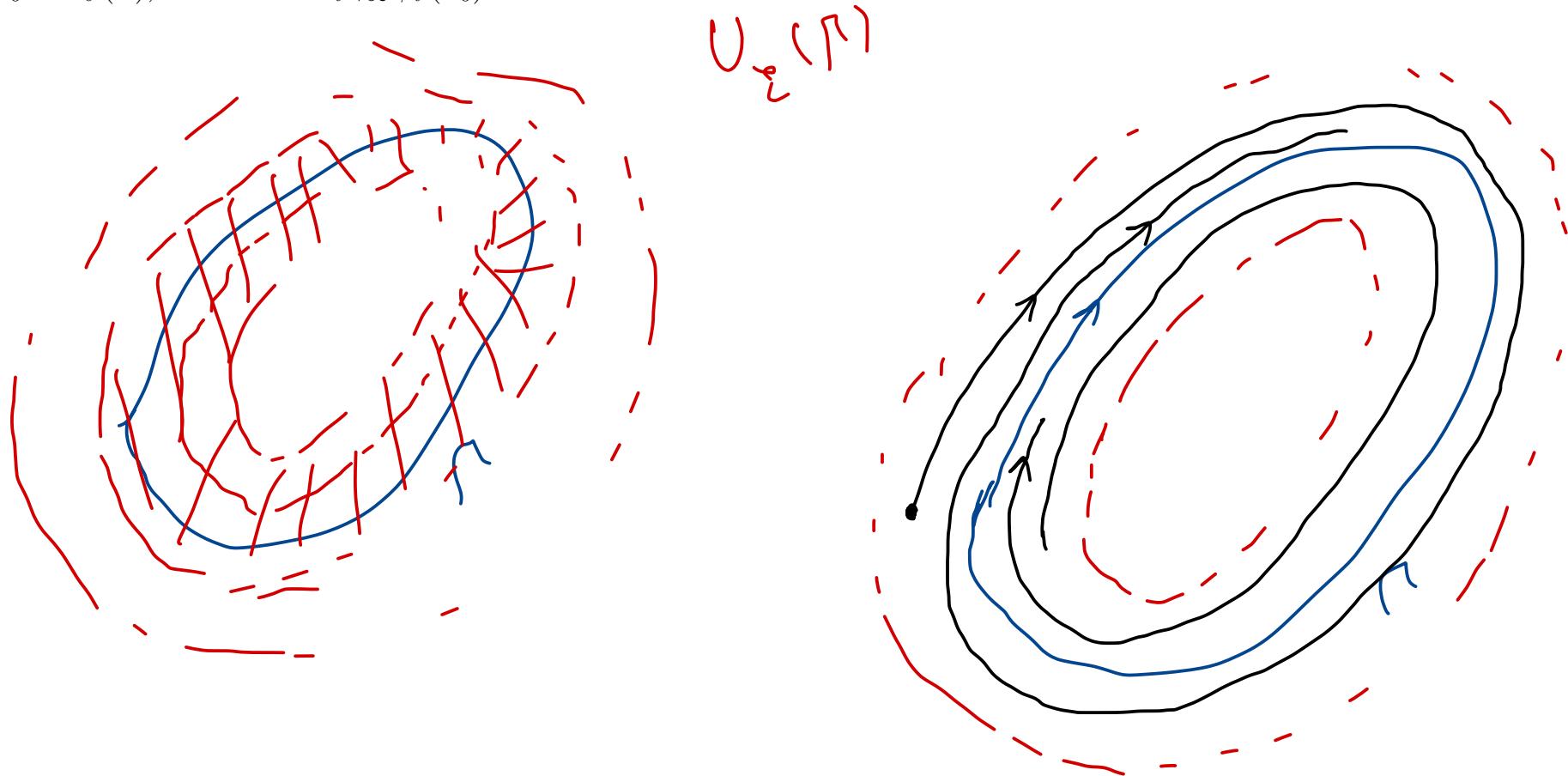
$\|x - \Gamma$

Analogous to equilibria, we can define the concept of stability for periodic orbits too. For arbitrary $x_0 \in \mathbb{R}^n$, let $\phi_t(x_0) = \phi(t, x_0)$ be the solution of the system with the initial condition $\phi(0, x_0) = x_0$.

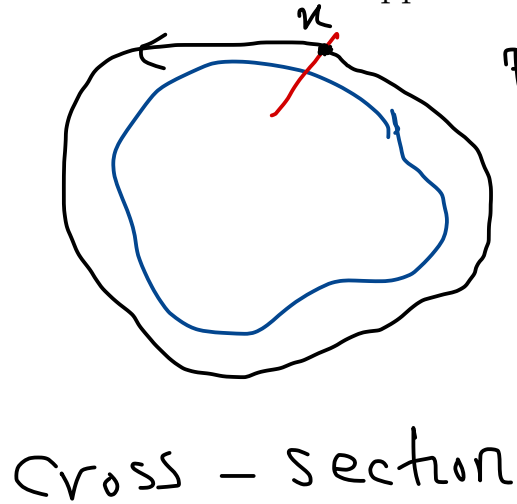
Definition 5.32. The cycle Γ is called stable if for every $\epsilon > 0$, there exists $\delta > 0$ such that for all $x_0 \in U_\delta(\Gamma)$ and for all $t \geq 0$, we have $\phi_t(x_0) \in U_\epsilon(\Gamma)$.

Definition 5.33. The cycle Γ is called unstable if it is not stable.

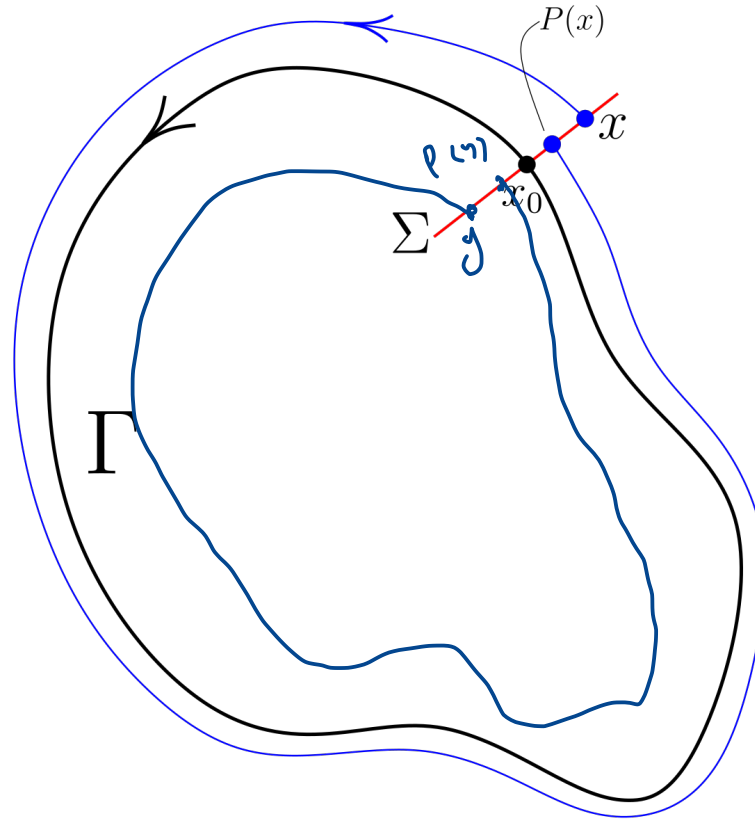
Definition 5.34. The cycle Γ is called asymptotically stable if it is stable, and there exists a $\delta > 0$ such that for all $x_0 \in U_\delta(\Gamma)$, we have $\lim_{t \rightarrow \infty} \phi_t(x_0) = \Gamma$.



- A standard approach to investigate the dynamics near a periodic orbit is to study an associated Poincaré map.



$$P(\xi) = \xi$$



$$x \in \Sigma \longrightarrow P(x) \in \Sigma$$

$$y \longmapsto P(y)$$

$$\xi_0 \longmapsto \xi_0$$

$$P(\xi_0) = \xi_0$$

$$P: \Sigma \longrightarrow \Sigma$$

$\subset \mathbb{R} \qquad \subset \mathbb{R}$

Figure 34: Here, Γ is a periodic orbit in the plane. The Poincaré map P maps $x \in \Sigma$ to $P(x) \in \Sigma$.

$$\xi \longrightarrow P(\xi) \longrightarrow P^2(\xi) = P(P(\xi)) \longmapsto \dots \qquad \xi_0$$

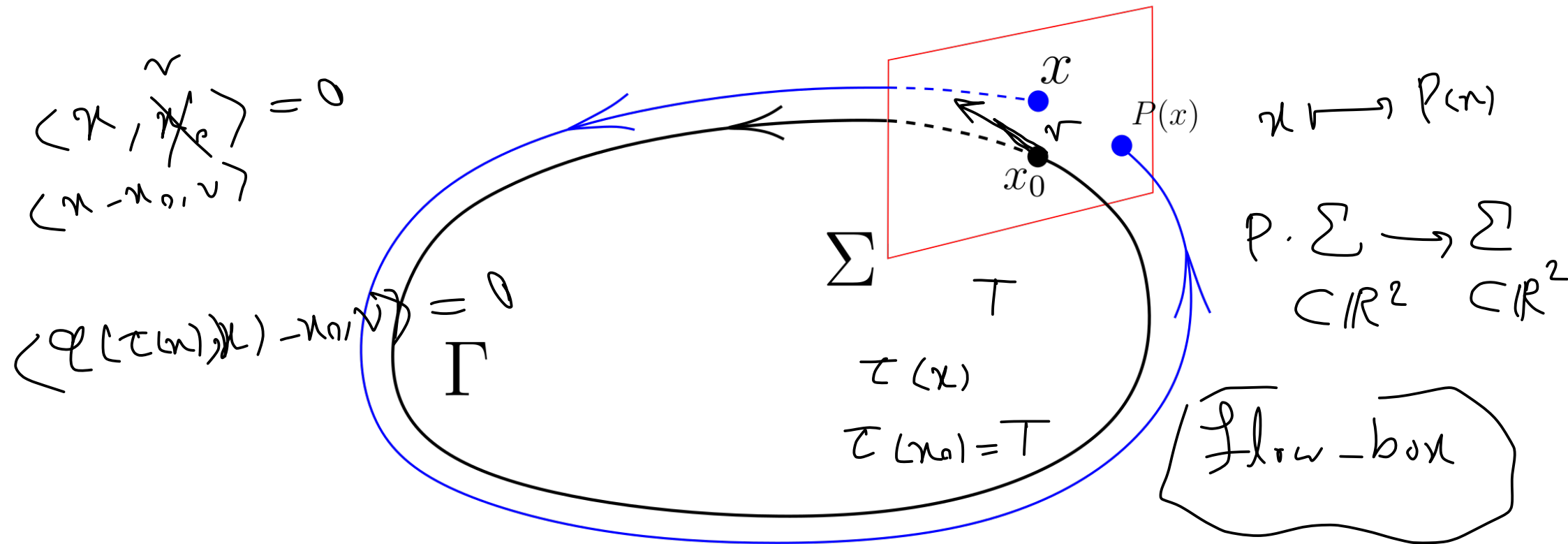


Figure 35: Here, Γ is a periodic orbit in the 3-dimensional space. The cross-section Σ is 2-dimensional. The Poincaré map P maps $x \in \Sigma$ to $P(x) \in \Sigma$.

5.5. Bifurcations

$$\dot{x} = f(x)$$

$$\dot{x} = f(x, \tau) \tag{5.32}$$

Consider a system

$$\dot{x} = f(x, \alpha),$$

where $x \in \mathbb{R}^n$ and $\alpha = (\alpha_1, \dots, \alpha_k) \in \mathbb{R}^k$. The variable x is the phase variable.

- We can think of system (5.32) as a model for a physical problem for which, α is a controlling parameter, such as temperature, pressure, etc.
- We can also think of system (5.32) as a family of systems. For each fixed α , we have a system of ODEs.

Example 5.35. Consider the system

$$\dot{x} = x^3 - \alpha x, \tag{5.33}$$

where $x \in \mathbb{R}$, and $\alpha \in \mathbb{R}$. For each fixed α , we have a system of ODEs. For instance,

(i) When $\alpha = 4$, we consider the system $\dot{x} = x^3 - 4x$.

(ii) When $\alpha = 0$, we consider the system $\dot{x} = x^3$.

(iii) When $\alpha = -10$, we consider the system $\dot{x} = x^3 + 10x$.

Example 5.36. Consider the system

$$\begin{aligned}\dot{x}_1 &= x_1 - \alpha_1 x_2^3 + \alpha_2 x_1 x_2, \\ \dot{x}_2 &= 4x_2 + \alpha_3 x_1 x_2^3 - \alpha_2 x_1^4,\end{aligned}\tag{5.34}$$

where $x = (x_1, x_2) \in \mathbb{R}^2$, and $\alpha = (\alpha_1, \alpha_2, \alpha_3) \in \mathbb{R}^3$. For each fixed α , we have a system of ODEs. For instance,

(i) When $\alpha = (0, 0, 0)$, we consider the system

$$\begin{aligned}\dot{x}_1 &= x_1 \\ \dot{x}_2 &= 4x_2.\end{aligned}\tag{5.35}$$

(ii) When $\alpha = (1, 2, 0)$, we consider the system

$$\begin{aligned}\dot{x}_1 &= x_1 - x_2^3 + 2x_1 x_2, \\ \dot{x}_2 &= 4x_2 - 2x_1^4,\end{aligned}\tag{5.36}$$

(iii) When $\alpha = (-1, 8, -2)$, we consider the system

$$\begin{aligned}\dot{x}_1 &= x_1 + x_2^3 + 8x_1 x_2, \\ \dot{x}_2 &= 4x_2 - 2x_1 x_2^3 - 8x_1^4,\end{aligned}\tag{5.37}$$

- In this course, we focus on the case that the parameter space is one dimensional, i.e. $\alpha \in \mathbb{R}$.

- Consider a system

$$\dot{x} = f(x, \alpha), \tag{5.38}$$

$$\alpha \in \mathbb{R}$$

where $x \in \mathbb{R}^n$ and $\alpha \in \mathbb{R}$. We say a bifurcation occurs at $\alpha = \alpha_c$ if the dynamics of the system changes suddenly at $\alpha = \alpha_c$.

- The parameter value α_c is called the bifurcation value.
- Recall from the first lecture that a mathematical model for understanding synchronization in networks is given by

$$\dot{x}_i = f_i(x_i) + \alpha \sum_{j=1}^N A_{ij} H_i(x_j - x_i), \quad \forall i \in \{1, \dots, N\}, \tag{5.39}$$

where $x_i \in \mathbb{R}^n$ ($n \geq 1$), $A = (A_{ij})$ is the adjacency matrix of the network, and $f_i, H_i \in \mathcal{C}^2(\mathbb{R}^n)$. As it was mentioned in the first lecture, there exists a parameter value α_c (and we also discussed for the linear case in Section 4.6), called critical coupling strength, such that for $\alpha < \alpha_c$, system (5.39) is not in synchrony but it gets into synchrony for $\alpha \geq \alpha_c$.

$$\alpha < \alpha_c$$

Example 5.37. Consider the system

$$\dot{x} = x^2 - \alpha, \quad x, \alpha \in \mathbb{R}. \tag{5.40}$$

$$x = x^2$$

$$x^2 - \alpha = 0$$

$$-\sqrt{\alpha}, \sqrt{\alpha}$$

- For $\alpha < 0$, system (5.40) has no equilibria. The phase portrait of this system is shown in Figure 36.

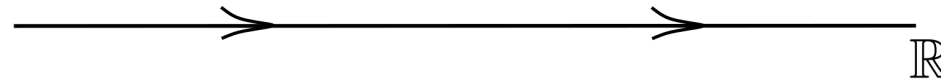


Figure 36: Case $\alpha < 0$.

- For $\alpha = 0$, system (5.40) has a nonhyperbolic equilibrium at $x = 0$. The phase portrait of this system is shown in Figure 37.

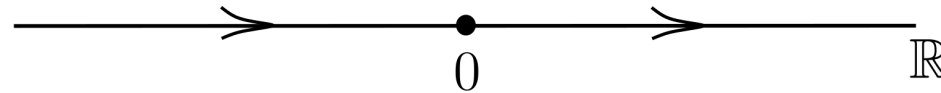


Figure 37: Case $\alpha = 0$.

- For $\alpha > 0$, system (5.40) has two hyperbolic equilibria at $x = -\sqrt{\alpha}$ (sink) and $x = \sqrt{\alpha}$ (source). The phase portrait of this system is shown in Figure 38.



Figure 38: Case $\alpha > 0$.

- A bifurcation occurs in this system at $\alpha = 0$.

$$x^2 - \alpha = 0$$

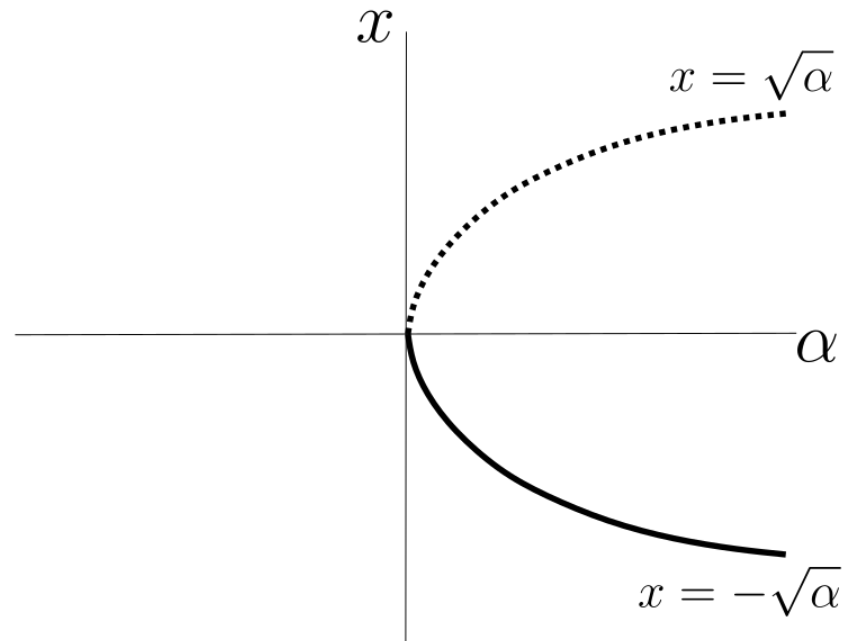


Figure 39: Bifurcation diagram for the system $\dot{x} = x^2 - \alpha$, where $x, \alpha \in \mathbb{R}$.

$$\lambda_c = 0$$

Hopf bifurcation

- Consider the nonlinear planar system

$$\begin{aligned} \dot{x} &= -y + x(\mu - x^2 - y^2), \\ \dot{y} &= x + y(\mu - x^2 - y^2), \end{aligned}$$

$$\begin{aligned} \dot{x} &= -y + \mu x \\ \dot{y} &= x + \mu y \end{aligned} \quad (5.41)$$

where $\mu \in \mathbb{R}$.

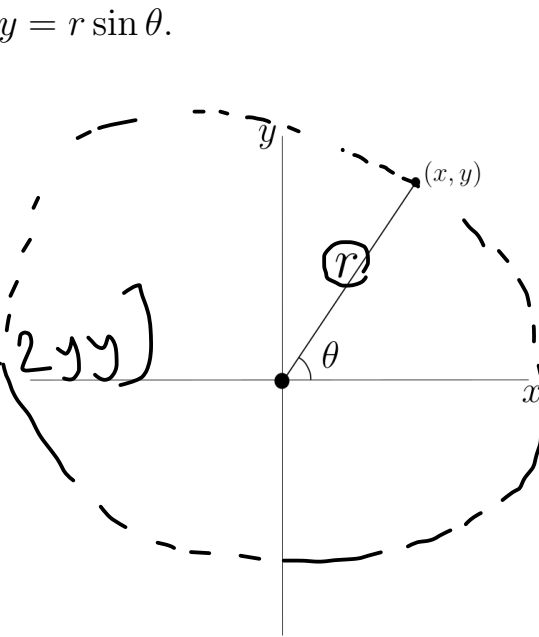
- The origin is an equilibrium. The linearization at the origin is $\begin{pmatrix} \mu & -1 \\ 1 & \mu \end{pmatrix}$. The corresponding eigenvalues are $\mu \pm i$.
- We can write system (5.41) in polar coordinates. Define

$$r := \sqrt{x^2 + y^2}, \quad \text{and} \quad \theta := \arctan \frac{y}{x}, \quad \text{where } (x \neq 0). \quad (5.42)$$

Equivalently, we have $x = r \cos \theta$ and $y = r \sin \theta$.

$$r = (x^2 + y^2)^{\frac{1}{2}}$$

$$\dot{r} = \frac{1}{2} (x^2 + y^2)^{-\frac{1}{2}} [2x\dot{x} + 2y\dot{y}]$$



$$\begin{aligned} \dot{r} \\ \dot{\theta} \end{aligned}$$

$$r = r_0$$

Figure 40: Polar coordinates.

- Writing system (5.41) in (r, θ) -coordinates, we have

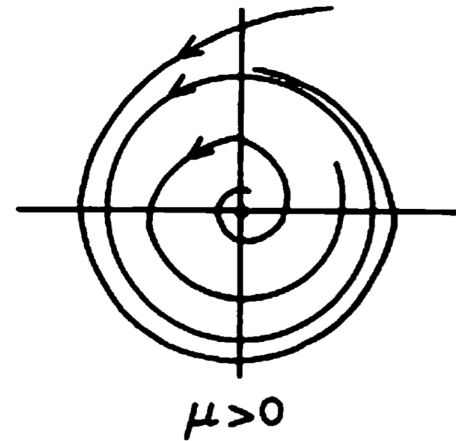
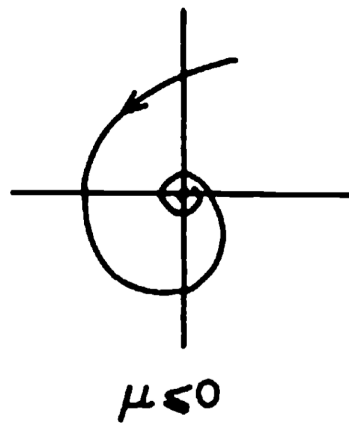
$$\begin{cases} \dot{r} = r(\mu - r^2), \\ \dot{\theta} = 1. \end{cases} \quad (5.43)$$

$\dot{r} = 0$

~~$r = 0$~~

~~$r = -\sqrt{\mu}, +\sqrt{\mu}$~~

- $\dot{r} = 0$ means that $r(t)$ is constant.
- Corresponding to the non-zero roots of \dot{r} , we have a periodic orbit for system (5.41), i.e. if \dot{r} vanishes at $r = r_0$ and $r_0 \neq 0$, then $(x(t), y(t))$, such that $[x(t)]^2 + [y(t)]^2 = r_0^2$, is a periodic orbit of system (5.41).
- Analyzing system (5.43), we can draw the phase portrait of system (5.41); see Figure 42.



$\mu > 0$

Figure 41: For $\mu \leq 0$, system (5.41) has an asymptotically stable equilibrium at the origin. Once μ becomes positive, the origin loses its stability and an asymptotically stable cycle gets born.

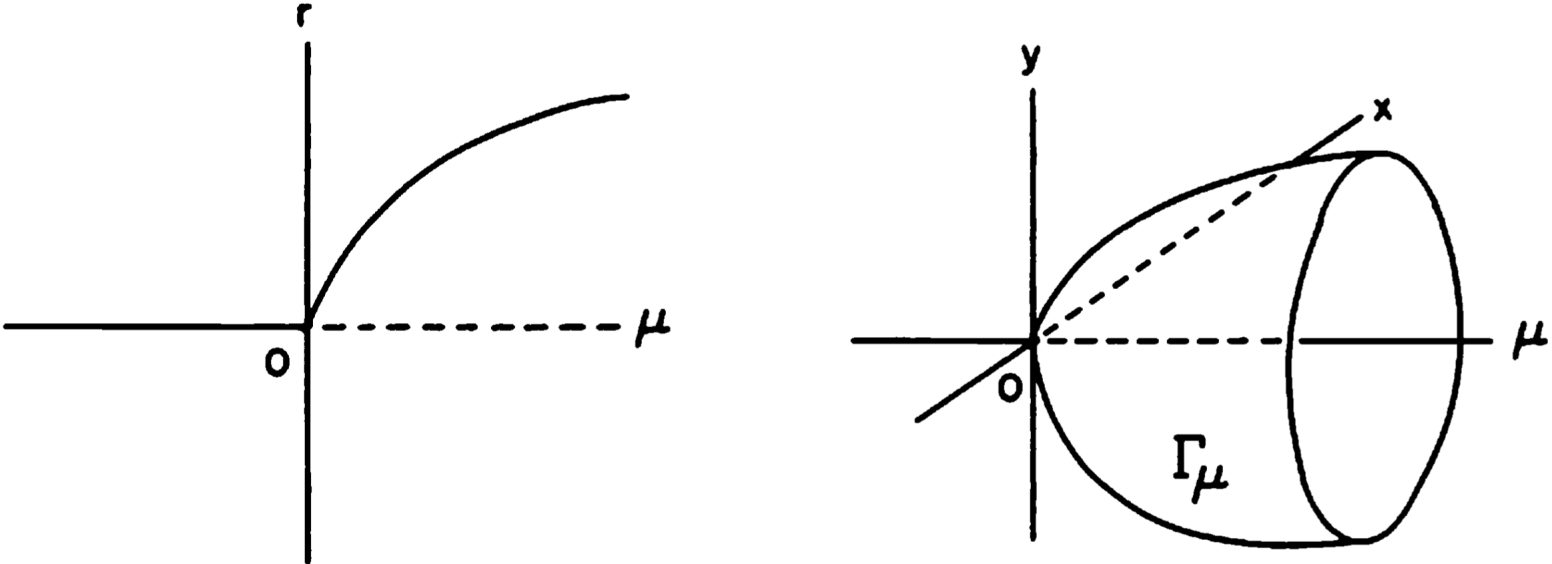


Figure 42: The periodic orbit associated with μ is shown by Γ_μ . This periodic orbit appears only when $\mu > 0$.

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